

For $\psi = Ae^{-br^2}$, normalize:

$$\int_0^\infty \int_0^{2\pi} \int_0^\pi A^2 e^{-2br^2} r^2 \sin\theta d\theta d\phi dr = 4\pi A^2 \int_0^\infty e^{-2br^2} r^2 dr = A^2 \left(\frac{\pi}{2b}\right)^{3/2} = 1$$

back cover of book

so $A = \left(\frac{2b}{\pi}\right)^{3/4}$. The \hat{H} operator is: $\hat{H} = -\frac{\hbar^2}{2m} \nabla^2 - \frac{e^2}{4\pi\epsilon_0 r}$

w/ $\nabla^2 \psi = \frac{2}{r} \psi' + \psi''$

$$\begin{aligned} \text{Then } \langle \psi | \nabla^2 \psi \rangle &= \int_0^\infty \int_0^{2\pi} \int_0^\pi \left(\frac{2b}{\pi}\right)^{3/2} e^{-br^2} \left[\frac{2}{r} (-2bre^{-br^2}) + (-2be^{-br^2} - 2br(-2bre^{-br^2})) \right] r^2 \sin\theta d\theta d\phi dr \\ &= 4\pi \left(\frac{2b}{\pi}\right)^{3/2} \int_0^\infty e^{-br^2} \left[-6be^{-br^2} + 4b^2 r^2 e^{-br^2} \right] r^2 dr \\ &= 4\pi \left(\frac{2b}{\pi}\right)^{3/2} \int_0^\infty \left[-6br^2 + 4b^2 r^4 \right] e^{-2br^2} dr \quad \leftarrow \text{back cover} \\ &= -4\pi \left(\frac{2b}{\pi}\right)^{3/2} \cdot \frac{3}{8} \left(\frac{\pi}{2b}\right)^{1/2} = -3b \end{aligned}$$

+ $\langle \psi | \hat{T} | \psi \rangle = \frac{3b\hbar^2}{2m}$
 $\leftarrow \equiv \frac{\hbar^2}{2m}$

so $\langle \psi | \frac{1}{r} | \psi \rangle = 4\pi \left(\frac{2b}{\pi}\right)^{3/2} \int_0^\infty \frac{1}{r} e^{-2br^2} r^2 dr = 4\pi \left(\frac{2b}{\pi}\right)^{3/2} \cdot \frac{1}{4b} = 2\sqrt{\frac{2b}{\pi}}$
 $\leftarrow \text{back cover}$

so $\langle \psi | U | \psi \rangle = -\frac{e^2}{4\pi\epsilon_0} \cdot 2\sqrt{\frac{2b}{\pi}} = -\frac{e^2 \sqrt{\frac{b}{\pi}}}{\epsilon_0 \pi^{3/4}}$
 $\leftarrow \equiv -\frac{e^2}{4\pi\epsilon_0 r}$

and the expectation value is: $\langle \psi | \hat{H} | \psi \rangle \equiv E_b = \frac{3b\hbar^2}{2m} - \frac{e^2 \sqrt{b}}{\sqrt{2}\epsilon_0 \pi^{3/4}}$

Now we minimize: $\frac{dE_b}{db} = \frac{3\hbar^2}{2m} - \frac{1}{2} \frac{e^2}{\sqrt{2}\epsilon_0 \pi^{3/4}} \frac{1}{\sqrt{b}} = 0$

or, solving for b : $\sqrt{b} = \frac{1}{2} \frac{e^2 \cdot 2m}{3\hbar^2 \sqrt{2}\epsilon_0 \pi^{3/4}} \Rightarrow b = \frac{m^2 e^4}{18\hbar^4 \epsilon_0^2 \pi^2}$

w/ $E_{\text{min}} = -\frac{me^4}{12\pi^2 \epsilon_0^2 \hbar^2} \equiv E_V$. The hydrogen ground state is $E_1 = -\frac{me^4}{32\pi^2 \epsilon_0^2 \hbar^2}$

so $E_V = \frac{8}{3\pi} E_1 \approx \frac{8}{3\pi} \approx .85$ so $E_V = -11.6 \text{ eV}$ (vs. -13.6 eV for hydrogen)

Problem 2 (8.4b)

For $\psi(x) = A x e^{-bx^2}$, we normalize: $\int_{-\infty}^{+\infty} A^2 x^2 e^{-2bx^2} dx = A^2 \cdot \frac{1}{4} \sqrt{\frac{\pi}{2b^3}} = 1 \Rightarrow A^2 = 4 \left(\frac{2b^3}{\pi}\right)^{1/2}$ back cover

The Hamiltonian is $\hat{H} = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + \frac{1}{2} m \omega^2 x^2$

For the kinetic piece: $\frac{d\psi}{dx} = A e^{-bx^2} - 2Abx e^{-bx^2}$
 $\frac{d^2\psi}{dx^2} = -2bx A e^{-bx^2} - 4Abx e^{-bx^2} + 4Ab^2 x^3 e^{-bx^2}$

then $\langle \psi | \hat{T} | \psi \rangle = -\frac{A^2 \hbar^2}{2m} \left[\int_{-\infty}^{+\infty} (-6bx^2 e^{-2bx^2} + 4b^2 x^4 e^{-2bx^2}) dx \right]$ (back cover)

$$= -\frac{A^2 \hbar^2}{2m} \left[-6b \cdot 2\sqrt{\pi} \cdot 2 \left(\frac{1}{2b}\right)^3 + 4b^2 \cdot 2\sqrt{\pi} \frac{4!}{2!} \left(\frac{1}{2b}\right)^5 \right]$$

$$= -\frac{\hbar^2}{2m} \cdot 4(2b^3)^{1/2} \left[-\frac{3b}{(2b)^{3/2}} + \frac{3b^2}{(2b)^{5/2}} \right] = -\frac{2\hbar^2}{m} \sqrt{2b} \left[-\frac{3}{2^{3/2}} + \frac{3}{2^{5/2}} \right]$$

$$= -\frac{\hbar^2}{m} \cdot 3b \left[-1 + \frac{1}{2} \right] = \frac{3b\hbar^2}{2m}$$

The potential piece is: $\langle \psi | U | \psi \rangle = A^2 \int_{-\infty}^{+\infty} \frac{1}{2} m \omega^2 x^4 e^{-2bx^2} dx$

$$= \frac{1}{2} m \omega^2 A^2 \left[\sqrt{\pi} \cdot \frac{4!}{2!} \cdot 2 \left(\frac{1}{2b}\right)^5 \right] = \frac{2m\omega^2}{\sqrt{\pi}} \sqrt{2b^3} \left[\sqrt{\pi} \frac{3}{4} \frac{1}{(2b)^{5/2}} \right]$$

$$= \frac{3m\omega^2}{2} \frac{\sqrt{2} b^{3/2}}{2^{5/2} b^{5/2}} = \frac{3m\omega^2}{8b}$$

and $E_b = \langle \psi | \hat{H} | \psi \rangle = \frac{3b\hbar^2}{2m} + \frac{3m\omega^2}{8b}$, now minimizing using b :

$$\frac{dE_b}{db} = \frac{3\hbar^2}{2m} - \frac{3m\omega^2}{8b^2} = 0 \Rightarrow b^2 = \frac{6m^2\omega^2}{24\hbar^2} \Rightarrow b = \frac{m\omega}{2\hbar}$$

the minimum is $E_{\min} = \frac{3\hbar^2}{2m} \left(\frac{m\omega}{2\hbar}\right) + \frac{3m\omega^2}{8} \cdot \frac{2\hbar}{m\omega} = \frac{3}{4}\hbar\omega + \frac{3}{4}\hbar\omega = \frac{3}{2}\hbar\omega$ ✓

Problem 3 (8.7)

For $\oplus \oplus \ominus$ we know $E = 4E_1$, ($4E_1 \approx -13.6 \text{ eV}$), so for a single electron, $E = 4E_1 \approx -54.4 \text{ eV}$, the difference between the ground state & this gives the ionization energy:

$$-54.4 \text{ eV} - (-79 \text{ eV}) = 24.6 \text{ eV}.$$

Problem 4

For the δ well, the single bound state is: $\psi_b = \frac{\sqrt{m\alpha}}{\hbar} e^{-m\alpha|x|/\hbar^2}$

The scattering states are $\psi_s = \begin{cases} A(e^{ikx} + \frac{i\beta}{1-i\beta} e^{-ikx}) & x \leq 0 \\ A \cdot \frac{1}{1-i\beta} e^{ikx} & x \geq 0 \end{cases}$ $\beta \equiv \frac{m\alpha}{\hbar^2 k}$, $k > 0$

we'll verify that $\langle \psi_b | \psi_s \rangle = 0$,

$$\langle \psi_b | \psi_s \rangle = \int_{-\infty}^{\infty} \psi_b^*(x) \psi_s(x) dx$$

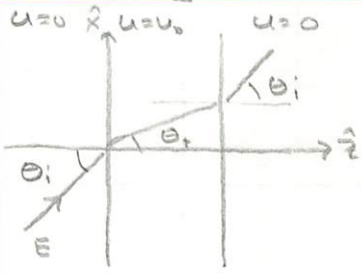
$$= \frac{A\sqrt{m\alpha}}{\hbar} \left[\int_{-\infty}^0 e^{x(i\hbar k + \beta)} dx + \frac{i\beta}{1-i\beta} \int_0^{\infty} e^{x(-i\hbar k + \beta)} dx + \frac{1}{1-i\beta} \int_0^{\infty} e^{x(i\hbar k - \beta)} dx \right]$$

$$= \frac{A\sqrt{m\alpha}}{\hbar} \left[\frac{1}{\hbar k(1-i\beta)} + \frac{i\beta}{(1-i\beta)(-i\hbar k + \beta)} - \frac{1}{(1-i\beta)(\hbar k(1+i\beta))} \right]$$

$$= \frac{A\sqrt{m\alpha}}{\hbar} \left[\frac{-i}{\hbar k(1-i\beta)} - \frac{\beta}{\hbar(1+\beta^2)} + \frac{i}{\hbar(1+\beta^2)} \right] = 0 \checkmark$$

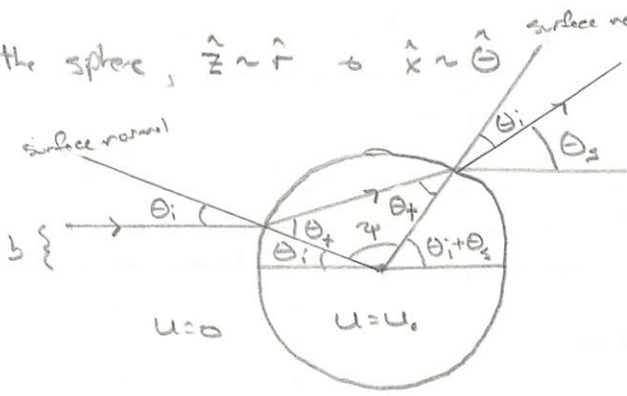
$$= \frac{-i(1+i\beta)}{\hbar(1+\beta^2)} = \frac{-i}{\hbar(1+\beta^2)} + \frac{\beta}{\hbar(1+\beta^2)}$$

Problem 5



$$\sin \theta_r = \frac{\sin \theta_i}{\sqrt{1 - u_0/E}}$$

For the sphere, $\hat{z} \sim \hat{r}$ to $\hat{x} \sim \hat{\theta}$



$$\psi + 2\theta_+ = \pi \Rightarrow \psi = \pi - 2\theta_+$$

$$\theta_i + \psi + (\theta_i + \theta_+) = \pi \Rightarrow \theta_s = 2(\theta_+ - \theta_i)$$

Now $\sin \theta_i = b/R$, to $\sin \theta_+ = \frac{\sin \theta_i}{\sqrt{1 - u_0/E}} = \frac{1}{\sqrt{1 - u_0/E}} \frac{b}{R}$

$$\theta_s = 2(\theta_+ - \theta_i) = 2(\sin^{-1}(\alpha b/R) - \sin^{-1}(b/R))$$