# Problem Set 9

Physics 322 Electrodynamics II

Due on Friday, April 5th, 2024

## Problem 1

Find the intensity of Thomson scattering for a charge q (initially at rest at the origin) that responds to an incident plane electromagnetic wave with electric field of the form:  $\mathbf{E} = E_0 \cos(k(z - ct))\hat{\mathbf{y}}$ . Write your result in terms of spherical coordinates, the "electron radius,"  $\ell$ , and the incident intensity  $I_0 = 1/2E_0^2/(\mu_0 c)$ . Why doesn't the intensity depend on the frequency  $\omega \equiv kc$ ?

## Problem 2

From the plane wave solutions to the wave equation (working in one spatial dimension here)

$$w(t) = w_0 e^{i2\pi(ft - kx)} \quad \text{with } f = kv, \tag{1}$$

we can make more general solutions by summing,

$$p(x,t) = \int_{-\infty}^{\infty} A(k)e^{i2\pi ft}e^{-i2\pi kx}dk.$$
(2)

where A(k) is a function of the wavenumber (an inverse length) k.

a. Given p(x,0) = u(x) as the initial data, find the expression for A(k) in terms of u(x) (and functions related to it – it might help to review the definition of the Fourier transform from the notes for January 31st, 2024).

**b.** Suppose the frequency f is itself a function of the wavenumber, f(k), so that

$$p(x,t) = \int_{-\infty}^{\infty} A(k) e^{i2\pi f(k)t} e^{-i2\pi kx} dk.$$
 (3)

If A(k) is sharply peaked about some  $k_0$ , then by Taylor expanding f about  $k_0$  and using that expression in (3), show that

$$p(x,t) \approx u(x - f'(k_0)t)e^{i\phi t}$$
(4)

where  $\phi$  is a constant. Up to a phase factor, then, p(x,t) is interpretable as a right-traveling wave that moves with speed  $f'(k_0)$  (the relevant group velocity here).

#### Problem 3

What is the energy required to build a spherical shell of radius R with total charge Q spread uniformly around its surface? How about a solid, uniformly charged sphere of radius R with that same charge Q? In each case, find the "classical charge radius."

#### Problem 4

Using the Thomson scattering setup, we can define the "scattering cross section" associated with electromagnetic radiation. For the incident wave:  $\mathbf{E}_{\rm in} = E_{\rm in} \cos(k(z-ct))\hat{\mathbf{n}}$ , the amount of power that goes through a disk of area  $d\sigma$  (with the disk's surface normal pointing in the direction of the incident wave,  $\hat{\mathbf{z}}$ ) is:  $S_{\rm in}d\sigma$ . The radiation that comes out (Thomson scattering) goes through a patch of a sphere of radius r, with  $d\mathbf{a} = r^2 \sin\theta d\theta d\phi \hat{\mathbf{r}}$ . The angular piece of the area patch defines the solid angle,  $d\Omega \equiv \sin\theta d\theta d\phi$ . Equating the incoming power with the outgoing power:

$$S_{\rm in}d\sigma = \mathbf{S}_{\rm out} \cdot \left(r^2 d\Omega \hat{\mathbf{r}}\right) \tag{5}$$

allows us to define the "differential scattering cross section," the function  $D(\theta) \equiv \frac{(r^2 \hat{\mathbf{r}} \cdot \mathbf{S}_{\text{out}})}{S_{\text{in}}}$ . Write out the expression for  $D(\theta)$  for Thomson scattering using the time-averaged Poynting vectors (the intensities), and averaging over the incident polarization.

## Problem 5

Griffiths 9.36 – transmission through multiple materials – feel free to use the matrix setup (or something similar to it) from problem 2 of Problem Set 8.

### Problem 6

Griffiths 9.38 - Scaring fish.